

# Violation of the action-reaction principle in the van der Waals interaction of excited atoms

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M. Donaire,  
arXiv:1604.07071 (2016)

M. Donaire,  
Phys. Rev. **A93**, 052706 (2016)

M. Donaire, R. Guerout, A. Lambrecht,  
Phys.Rev.Lett. **115**, 033201 (2015)

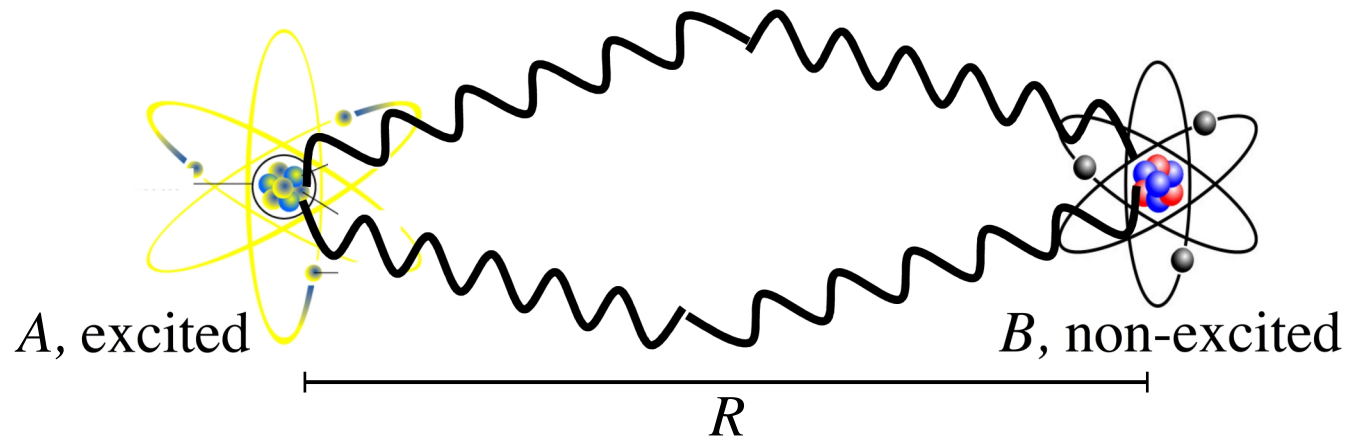


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# Outline



## 0. Two-atom phase shift $\neq$ single atom phase shifts $\neq$ vdW potentials $\neq$ photon frequency shift

P.R. Berman,  
Phys. Rev. A **91**, 042127 (2015)

P.Milonni & S.Rafsanjani,  
Phys. Rev. A **92**, 062711 (2015)

M.Donaire, R.Guerout, A.Lambrecht,  
Phys.Rev.Lett. **115**, 033201 (2015)

M. Donaire,  
Phys. Rev. A **93**, 052706 (2016)

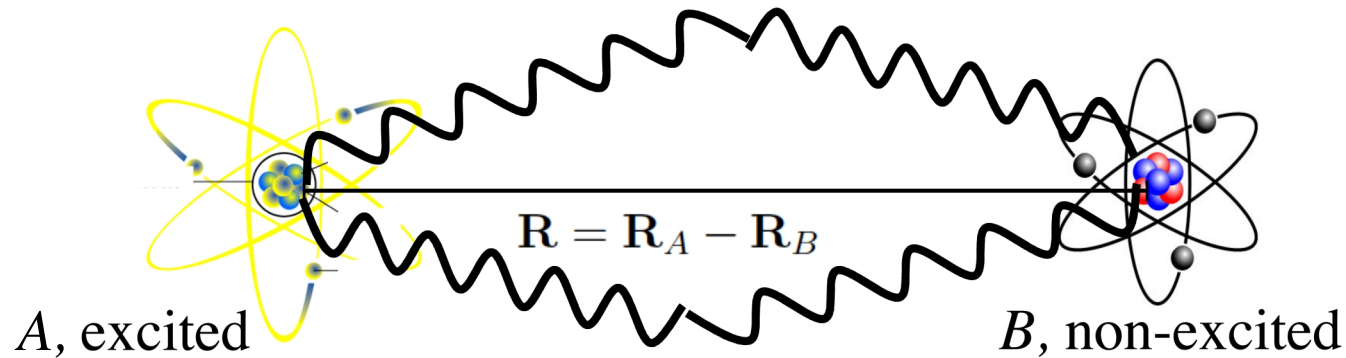
## i. Computation of the dynamical vdW forces

M. Donaire,  
Phys. Rev. A **93**, 052706 (2016)

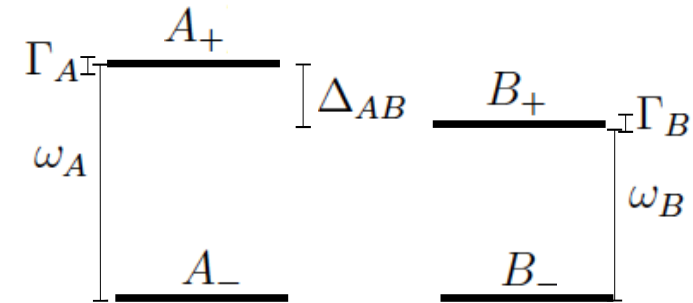
## ii. Violation of the action-reaction principle

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## i- Computation of the dynamical vdW forces



$$H_0 = \hbar\omega_A|A_+\rangle\langle A_+| + \hbar\omega_B|B_+\rangle\langle B_+| + \sum_{\mathbf{k},\epsilon} \hbar\omega(a_{\mathbf{k},\epsilon}^\dagger a_{\mathbf{k},\epsilon} + 1/2)$$



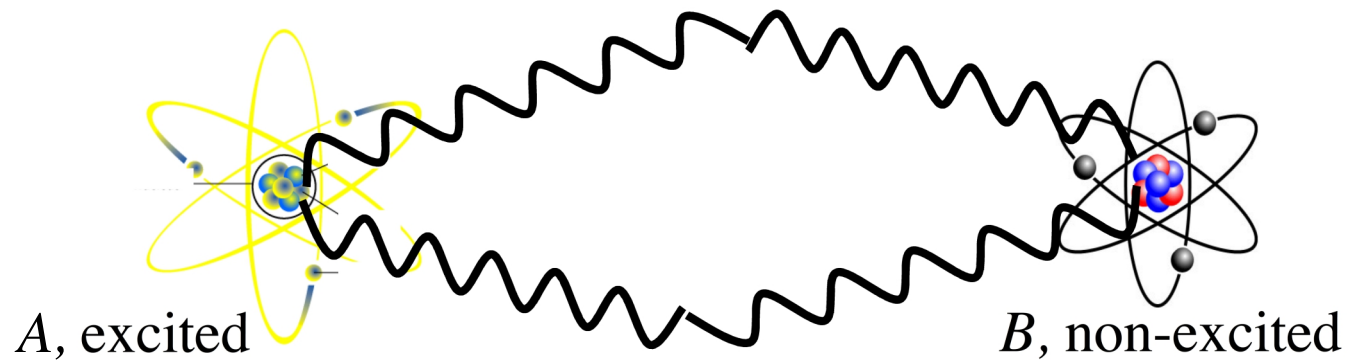
$$W = W_A + W_B = -\mathbf{d}_A \cdot \mathbf{E}(\mathbf{R}_A) - \mathbf{d}_B \cdot \mathbf{E}(\mathbf{R}_B)$$

$$\mathbf{E}(\mathbf{R}_{A,B}) = i \sum_{\mathbf{k},\epsilon} \sqrt{\frac{\hbar ck}{2\mathcal{V}\epsilon_0}} [\epsilon a_{\mathbf{k}} e^{i\mathbf{k} \cdot \mathbf{R}_{A,B}} - \epsilon^* a_{\mathbf{k}}^\dagger e^{-i\mathbf{k} \cdot \mathbf{R}_{A,B}}]$$

$$\Delta_{AB} \equiv \omega_A - \omega_B$$

$$\Gamma_A \ll |\Delta_{AB}|$$

## i- Computation of the dynamical vdW forces



Why dynamical?

1. Finite life-times  $\sim 1/\Gamma_A$
2. Periodic excitation transfer with probability  $\sim \langle W \rangle / \hbar \Delta_{AB} \ll 1$

## i- Computation of the dynamical vdW forces – Time-dependent pert. Th.

Wave-function formalism in the Schrödinger picture  $|\Psi(T)\rangle = \mathbb{U}(T, 0)|\Psi(0)\rangle$

$$\langle \mathcal{O}(T) \rangle = \langle \Psi(T) | \mathcal{O} | \Psi(T) \rangle = \langle \Psi(0) | \mathbb{U}^\dagger(T, 0) \mathcal{O} \mathbb{U}(T, 0) | \Psi(0) \rangle$$

$$\mathbb{U}(T, 0) = \mathcal{T} \exp \left\{ -i\hbar^{-1} \int_0^T dt [H_0 + W] \right\}$$

## i- Computation of the dynamical vdW forces – Time-dependent pert. Th.

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$$\mathbb{U}(T, 0) = \mathcal{T} \exp \left\{ -i\hbar^{-1} \int_0^T dt [H_0 + W] \right\}$$

$$\mathbb{U}(T, 0) = \mathbb{U}_0(T, 0) + \delta\mathbb{U}(T, 0)$$

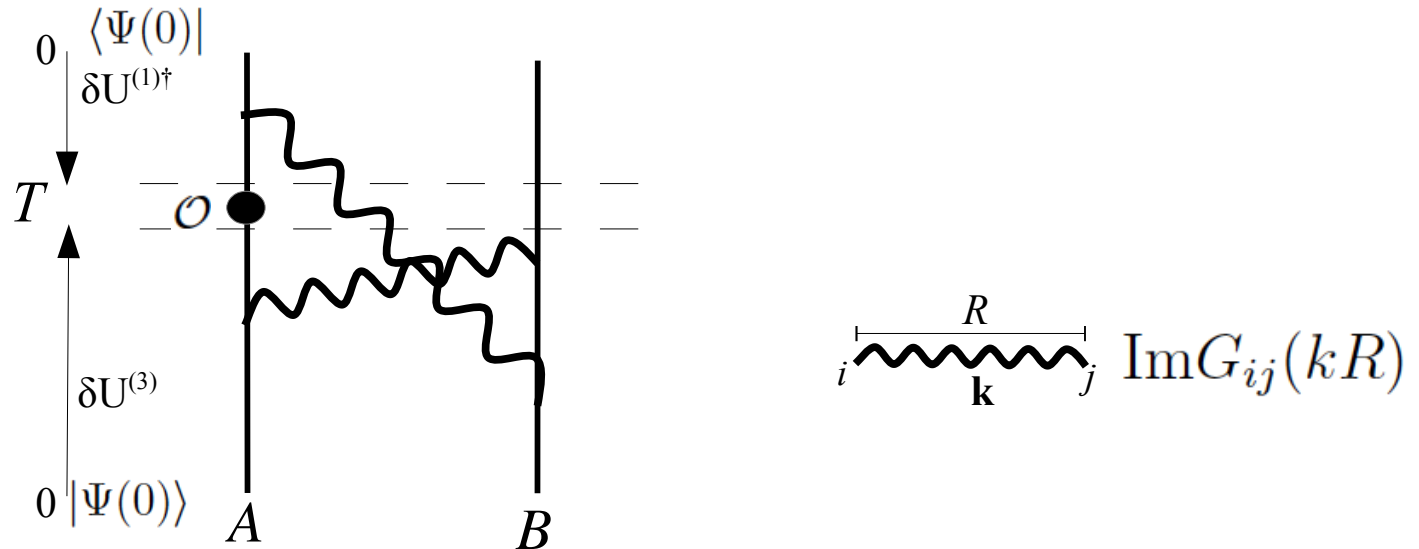
$$\mathbb{U}_0(T, 0) = \exp \left\{ -i\hbar^{-1} H_0 T \right\}$$

$$\begin{aligned} \delta\mathbb{U}(T) \simeq & (-i/\hbar) \int_0^T dt \mathbb{U}_0(T-t) W \mathbb{U}_0(t) \xleftarrow{\delta\mathbb{U}^{(1)}(T)} \\ & + (-i/\hbar)^2 \int_0^T dt \int_0^t dt' \mathbb{U}_0(T-t) W \mathbb{U}_0(t-t') W \mathbb{U}_0(t') \xleftarrow{\delta\mathbb{U}^{(2)}(T)} \\ & + (-i/\hbar)^3 \int_0^T dt \int_0^t dt' \int_0^{t'} dt'' \mathbb{U}_0(T-t) W \mathbb{U}_0(t-t') W \mathbb{U}_0(t'-t'') W \mathbb{U}_0(t'') \xleftarrow{\delta\mathbb{U}^{(3)}(T)} \end{aligned}$$

## i- Computation of the dynamical vdW forces – Time-dependent pert. Th.

Wave-function formalism in the Schrödinger picture  $|\Psi(T)\rangle = \mathbb{U}(T, 0)|\Psi(0)\rangle$

$$\langle \mathcal{O}(T) \rangle = \langle \Psi(T) | \mathcal{O} | \Psi(T) \rangle = \langle \Psi(0) | \mathbb{U}^\dagger(T, 0) \mathcal{O} \mathbb{U}(T, 0) | \Psi(0) \rangle$$



$$\begin{aligned} \delta \mathbb{U}(T) &\simeq (-i/\hbar) \int_0^T dt \mathbb{U}_0(T-t) W \mathbb{U}_0(t) \longleftarrow \delta \mathbb{U}^{(1)}(T) \\ &+ (-i/\hbar)^2 \int_0^T dt \int_0^t dt' \mathbb{U}_0(T-t) W \mathbb{U}_0(t-t') W \mathbb{U}_0(t') \longleftarrow \delta \mathbb{U}^{(2)}(T) \\ &+ (-i/\hbar)^3 \int_0^T dt \int_0^t dt' \int_0^{t'} dt'' \mathbb{U}_0(T-t) W \mathbb{U}_0(t-t') W \mathbb{U}_0(t'-t'') W \mathbb{U}_0(t'') \\ &\longleftarrow \delta \mathbb{U}^{(3)}(T) \end{aligned}$$

## i- Computation of the dynamical vdW forces

Van der Waals force and potential on atom  $A$

$$\begin{aligned}\langle \mathbf{F}_A(T) \rangle &= \partial_T \langle \mathbf{P}_A(T) \rangle = -i\hbar \partial_T \langle \Psi(0) | \mathbb{U}^\dagger(T) \nabla_{\mathbf{R}_A} \mathbb{U}(T) | \Psi(0) \rangle + O(v/c) \\ &= -\langle \nabla_{\mathbf{R}_A} W_A(T, \mathbf{R}_A) \rangle\end{aligned}$$

Idem for atom B, vdW potential on atom B

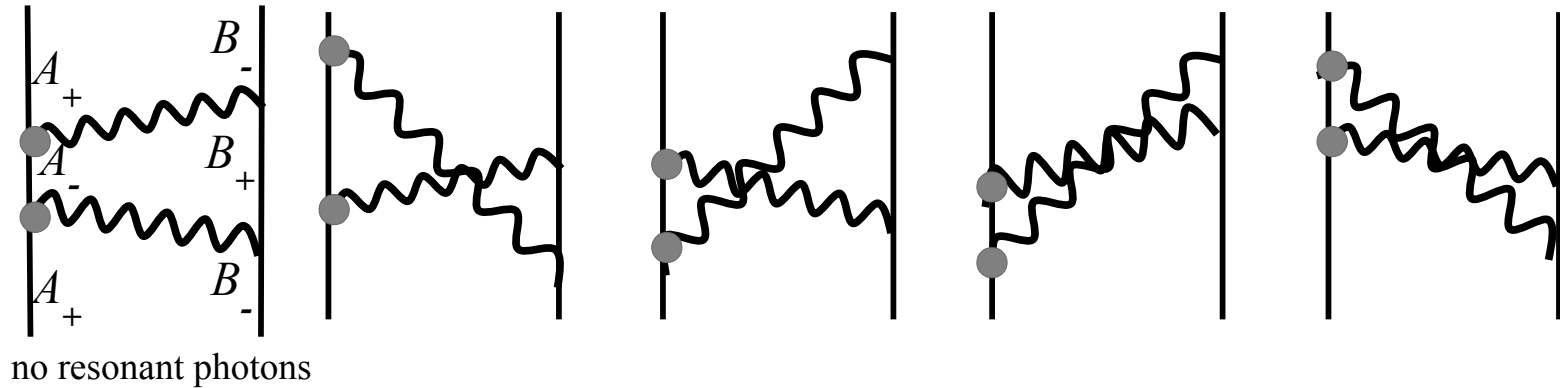
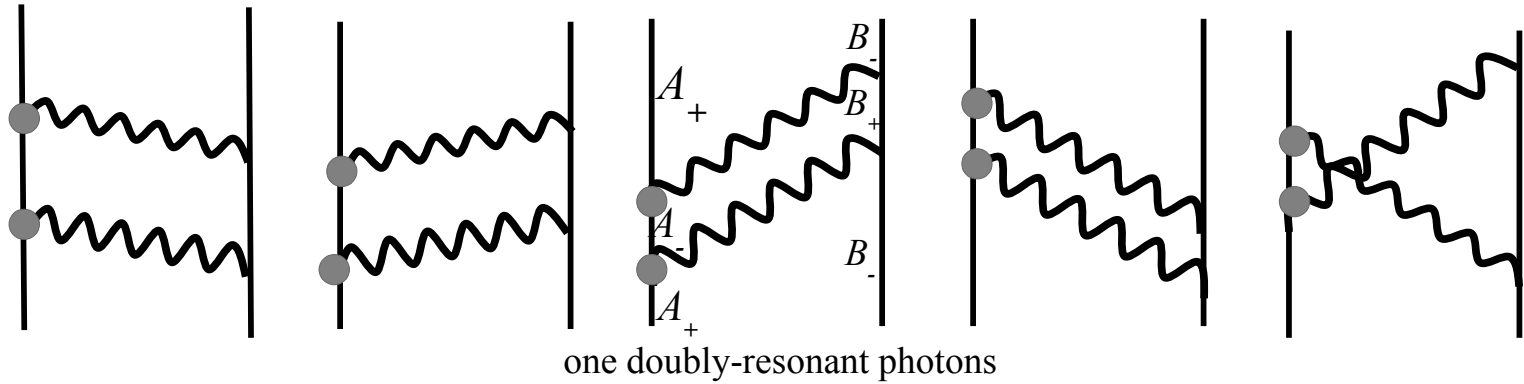
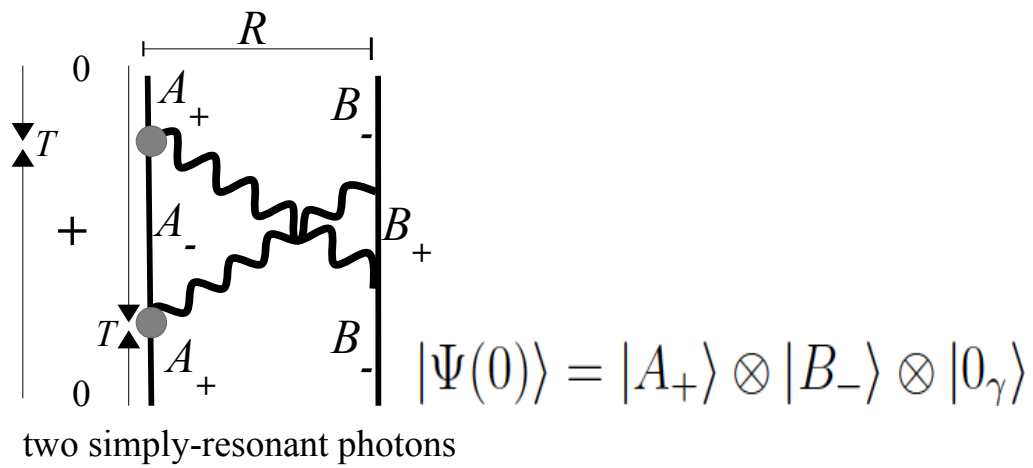
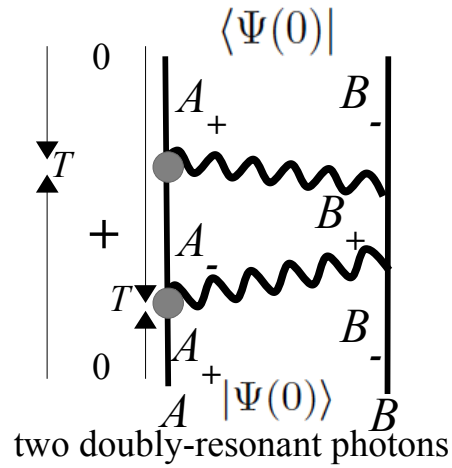
$$\langle \mathbf{F}_B(T) \rangle = -\langle \nabla_{\mathbf{R}_B} W_B(T, \mathbf{R}_B) \rangle$$



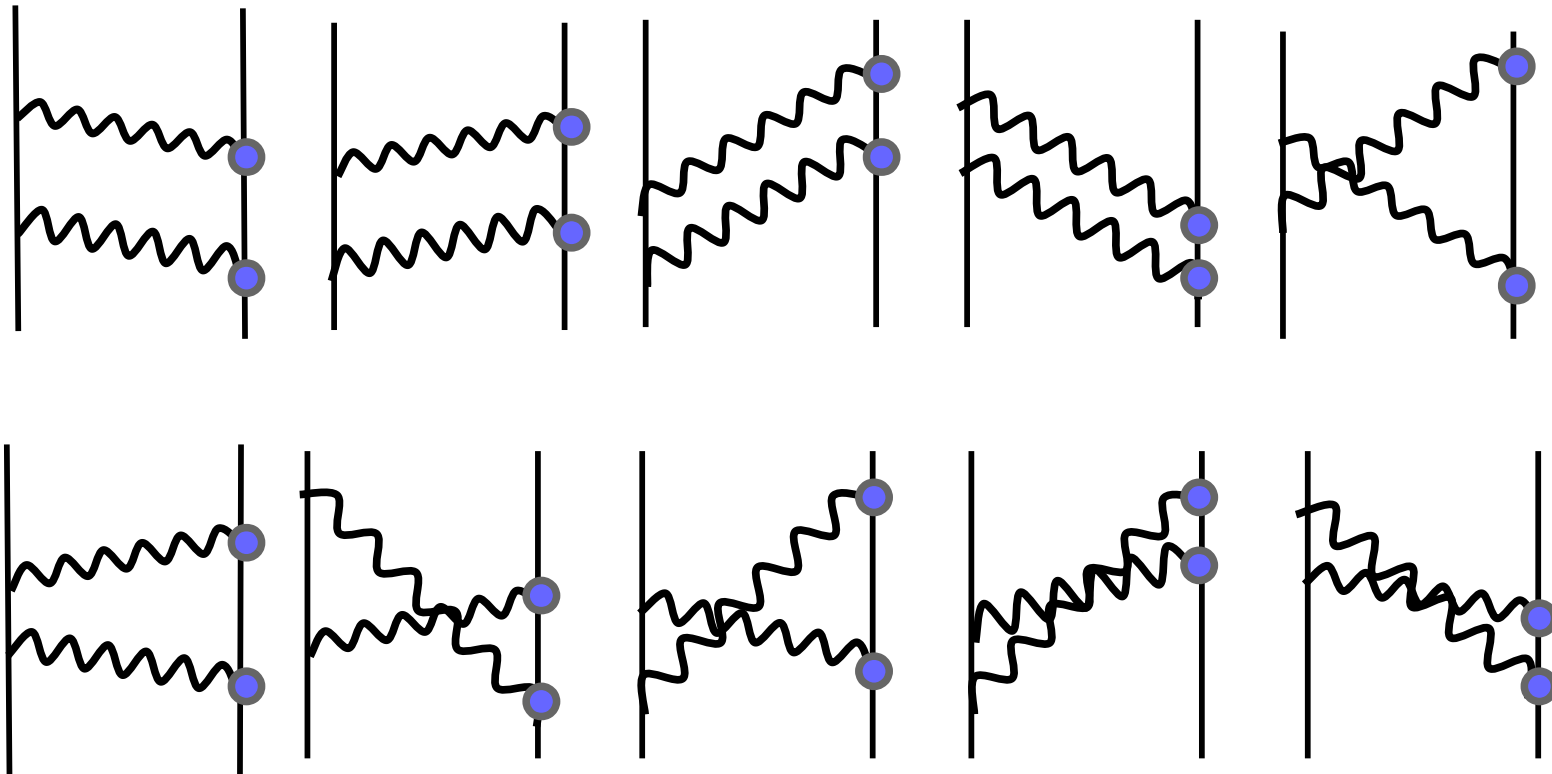
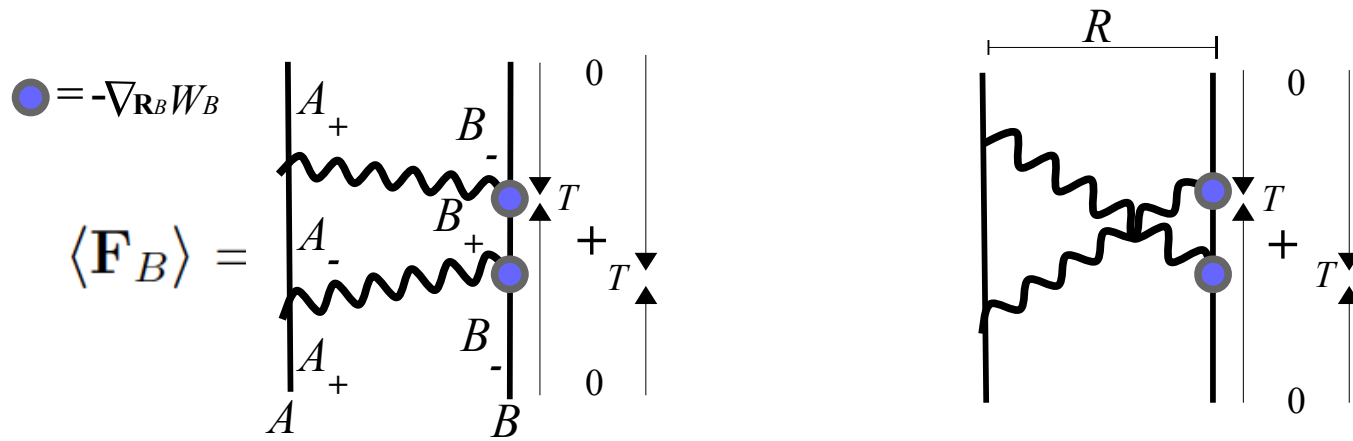
# i- Computation of the dynamical vdW forces

● =  $-\nabla_{\mathbf{R}_A} W_A$

$\langle \mathbf{F}_A \rangle =$



# i- Computation of the dynamical vdW forces



# i- Computation of the dynamical vdW forces

## Relevant diagrams for action-reaction violation

1. Non-resonant photons mediate the reciprocal exchange of momentum between the atoms, and generate the usual London and Casimir-Polder forces,

whereas

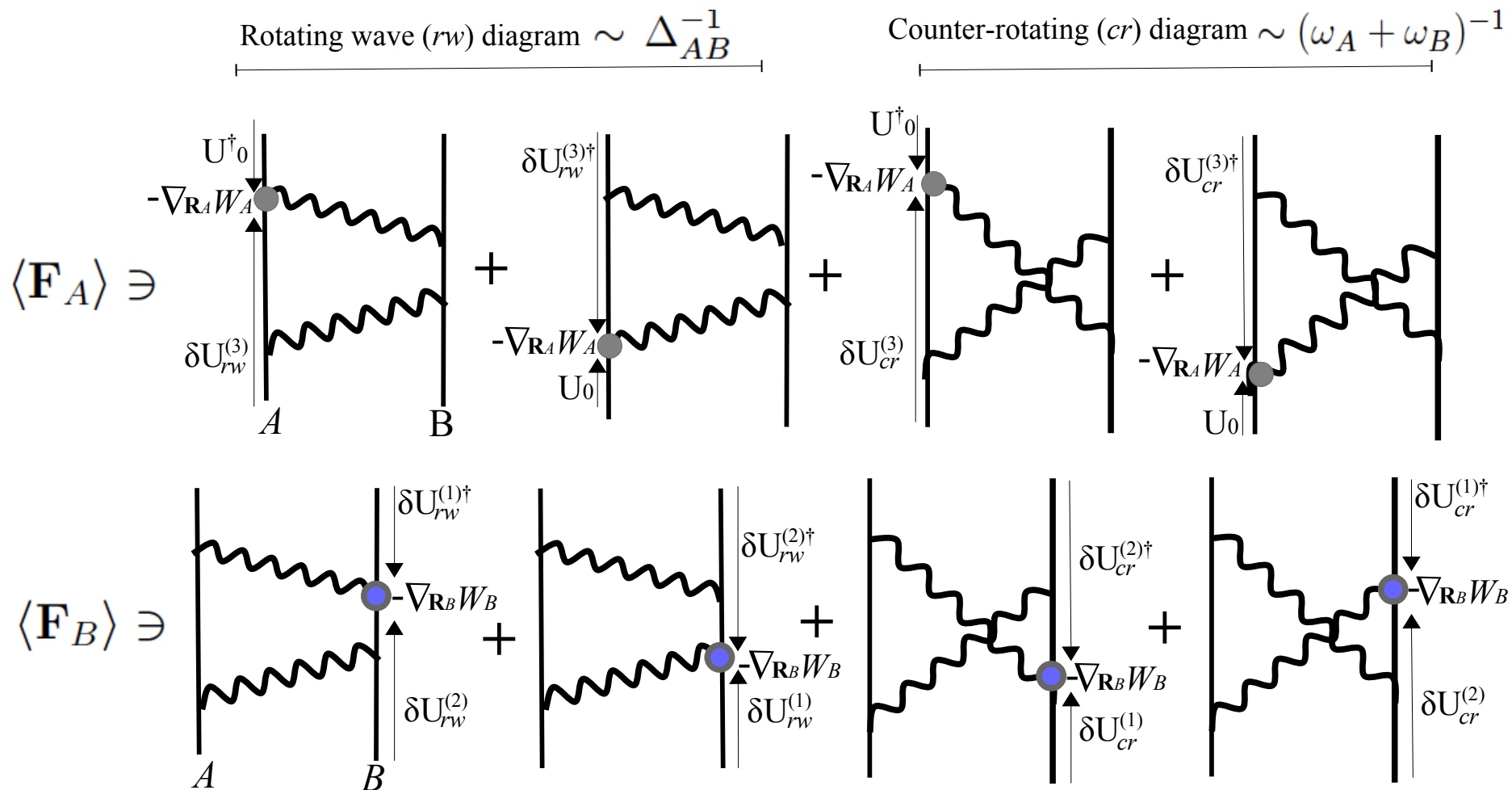
2. resonant photons mediate the transfer of the excitation between the atoms, and 'may' carry momentum off the two-atom system—see later—

1.&2. => Only diagrams with resonant photons may lead to action-reaction violation

# i- Computation of the dynamical vdW forces

## Relevant diagrams for action-reaction violation

3. Under realistic conditions (adiabatic excitation), only the following diagrams survive



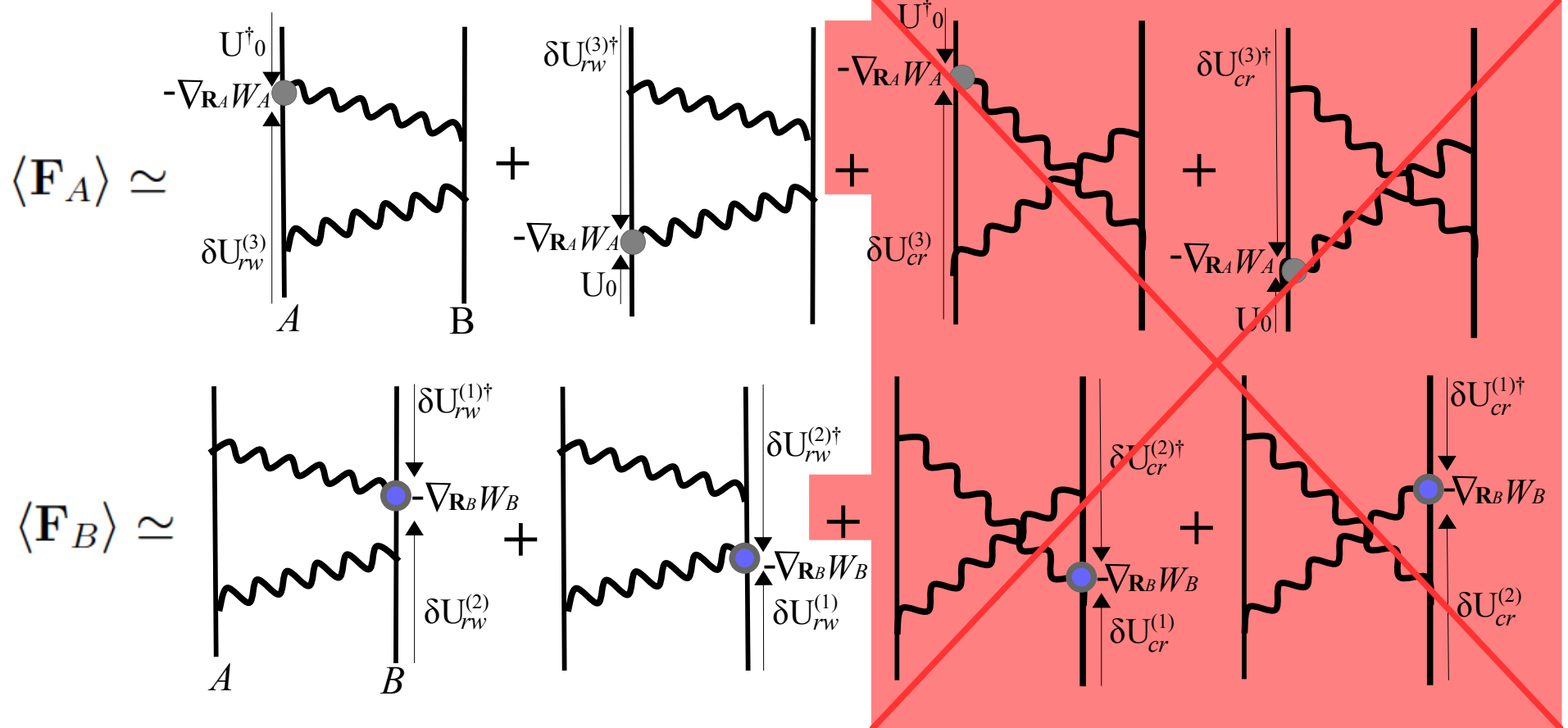
# i- Computation of the dynamical vdW forces

Quasiresonant approximation ~ Rotating wave approximation

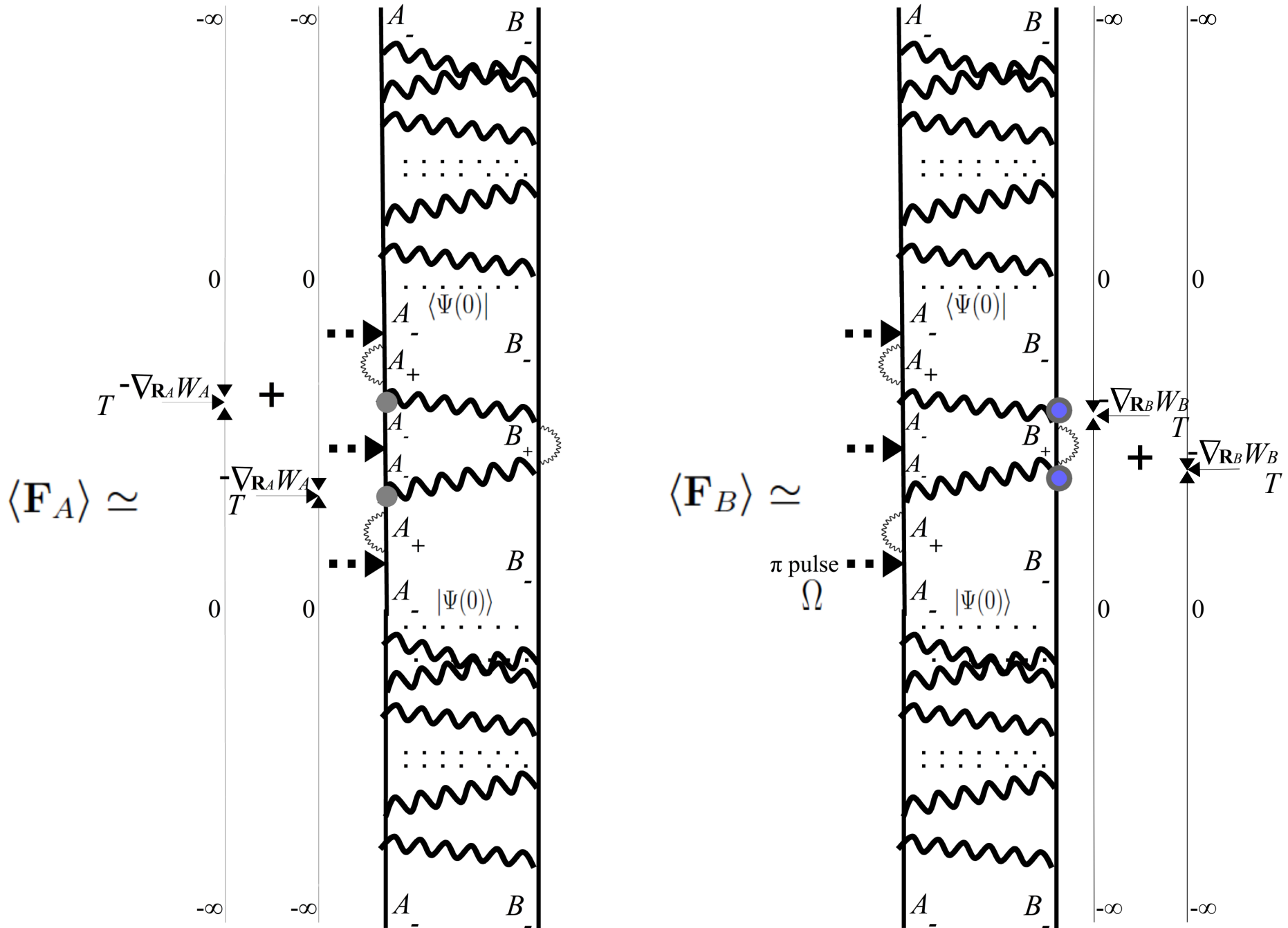
$$\Delta_{AB} \ll \omega_A, \omega_B$$

Rotating wave (*rw*) diagram  $\sim \Delta_{AB}^{-1}$

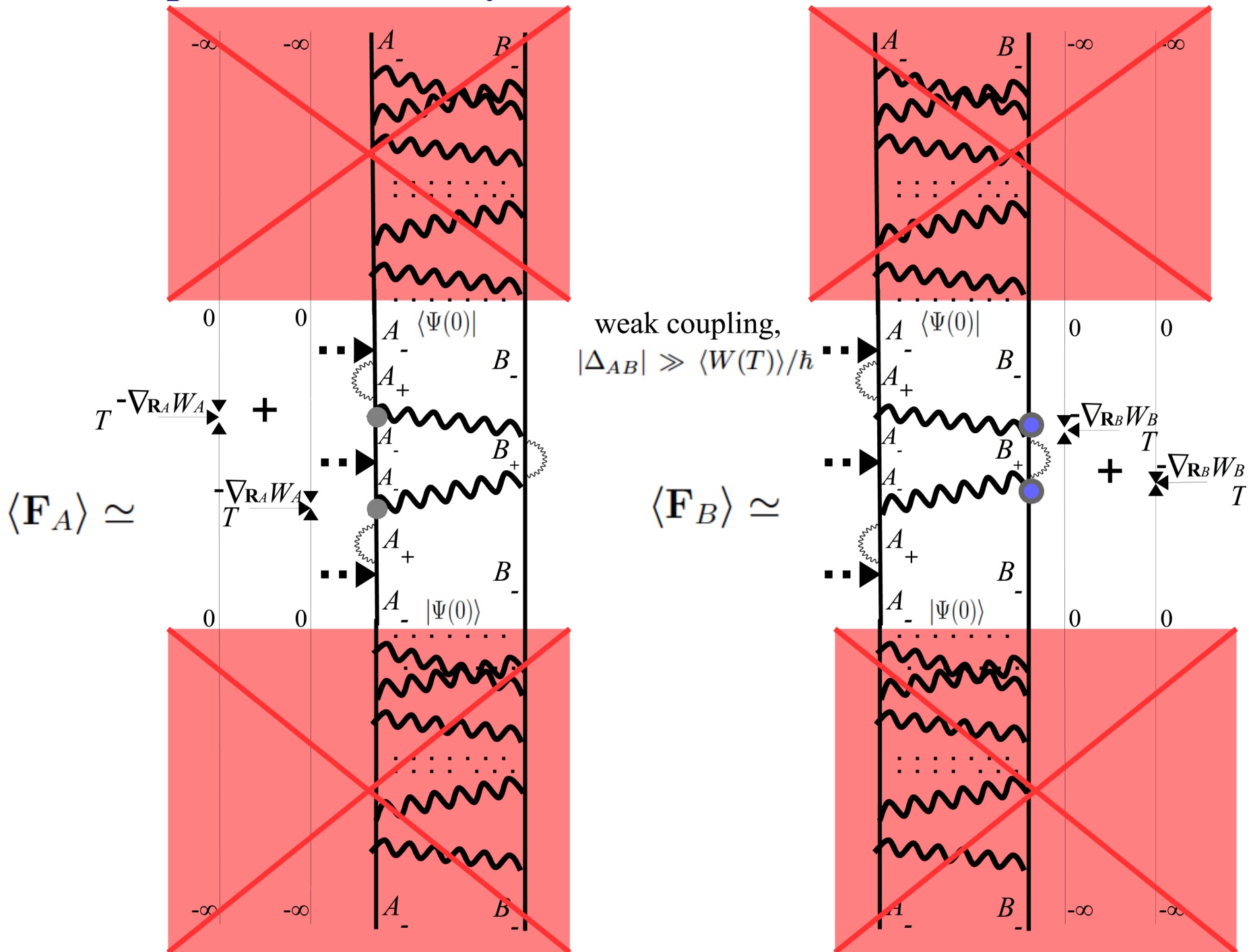
Counter-rotating (*cr*) diagram  $\sim (\omega_A + \omega_B)^{-1}$



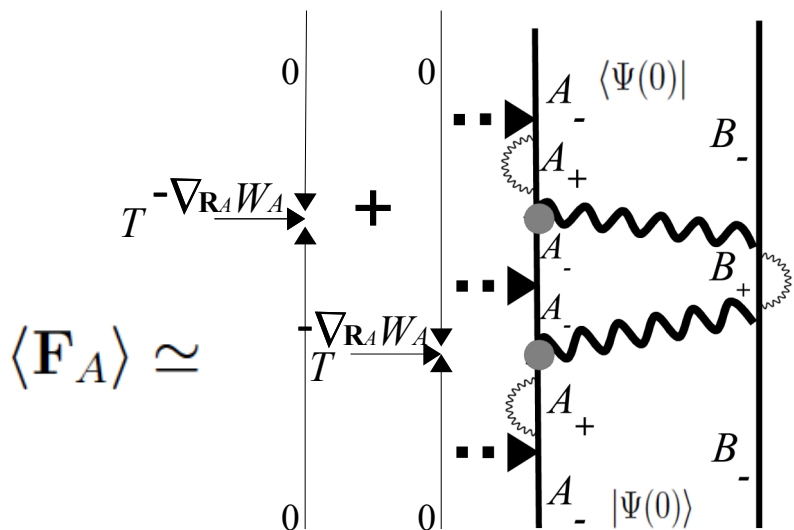
# i- Computation of the dynamical vdW forces



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## i- Computation of the dynamical vdW forces



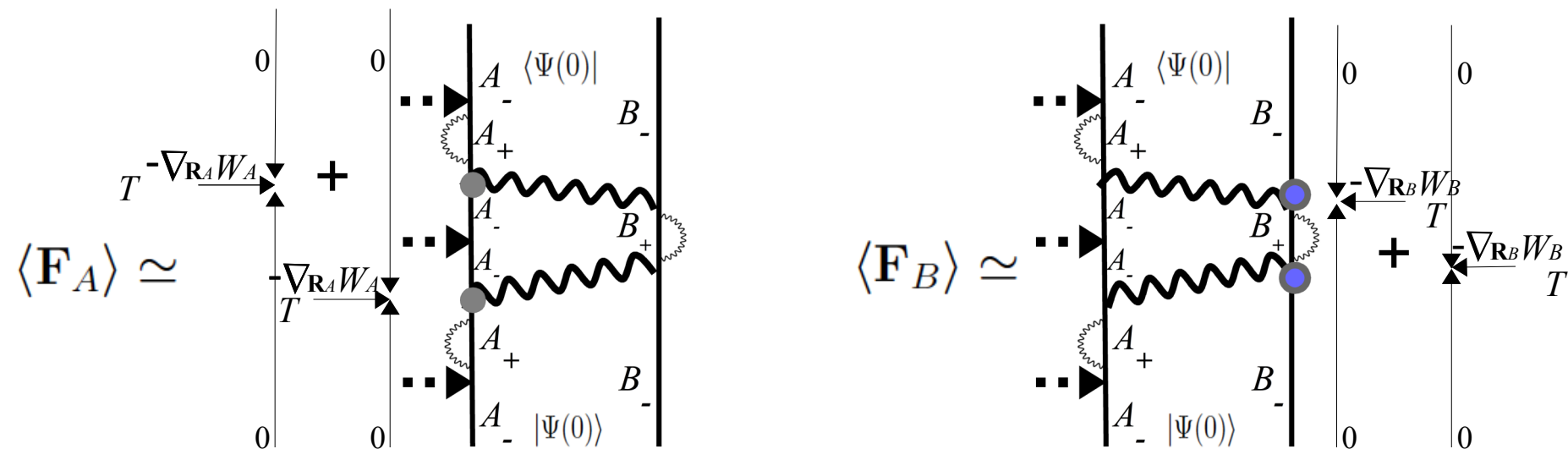
$$\begin{aligned} \langle \mathbf{F}_A(T) \rangle \simeq & \mathcal{U}^{ijpq} \left\{ k_A^4 e^{-\Gamma_A T} \nabla_{\mathbf{R}} \left[ \text{Re } G_{ij}^{(0)}(\mathbf{R}, \omega_A) \text{Re } G_{pq}^{(0)}(\mathbf{R}, \omega_A) - \text{Im } G_{ij}^{(0)}(\mathbf{R}, \omega_A) \text{Im } G_{pq}^{(0)}(\mathbf{R}, \omega_A) \right] \right. \\ & + \frac{k_B^4 \Omega^2 e^{-(\Gamma_A + \Gamma_B)T/2}}{2(\Delta_{AB}^2 - \Omega^2)} \nabla_{\mathbf{R}} \left[ \left( \text{Re } G_{ij}^{(0)}(\mathbf{R}, \omega_B) \text{Re } G_{pq}^{(0)}(\mathbf{R}, \omega_B) - \text{Im } G_{ij}^{(0)}(\mathbf{R}, \omega_B) \text{Im } G_{pq}^{(0)}(\mathbf{R}, \omega_B) \right) \right. \\ & \times (\cos(\Delta_{AB}T) + \cos[\Delta_{AB}(T - \pi/\Omega)]) - 2 \text{Re } G_{ij}^{(0)}(\mathbf{R}, \omega_B) \text{Im } G_{pq}^{(0)}(\mathbf{R}, \omega_B) \\ & \left. \left. \times (\sin(\Delta_{AB}T) + \sin[\Delta_{AB}(T - \pi/\Omega)]) \right] \right\} \end{aligned}$$

$$\mathbb{G}^{(0)}(kR) = \frac{k e^{ikR}}{-4\pi} [\gamma/kR + i\beta/(kR)^2 - \beta/(kR)^3] \quad \gamma = \mathbb{I} - \mathbf{R}\mathbf{R}/R^2, \beta = \mathbb{I} - 3\mathbf{R}\mathbf{R}/R^2$$

$$\mathcal{U}^{ijpq} = d_A^i d_B^j d_B^p d_A^q / \epsilon_0^2 \hbar \Delta_{AB}$$



# i- Computation of the dynamical vdW forces

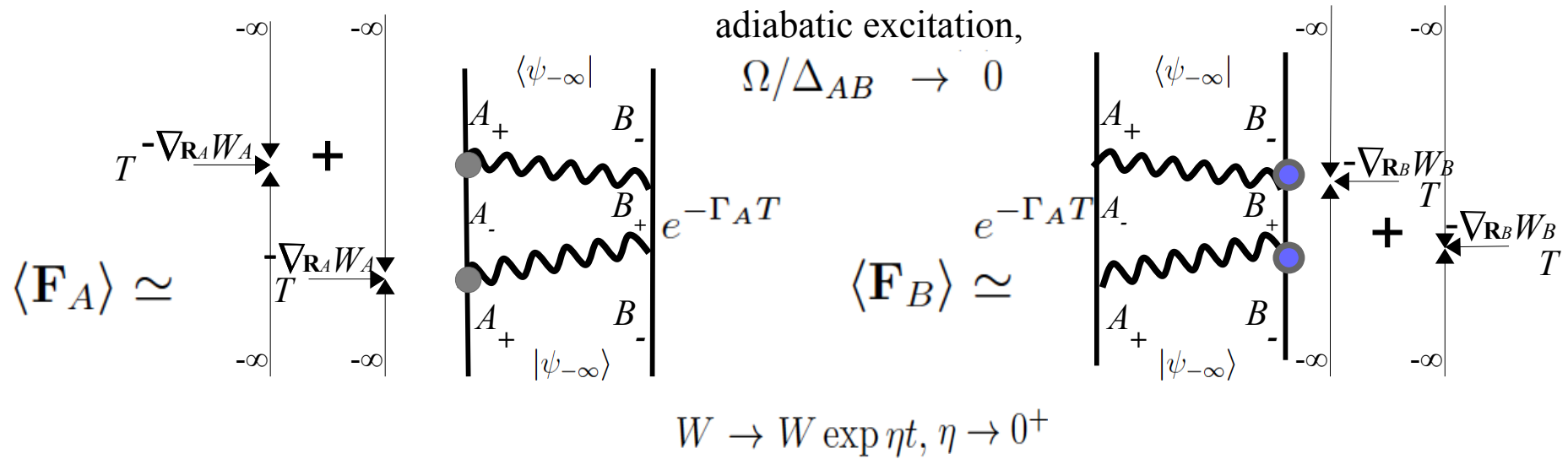


$$\begin{aligned} \langle \mathbf{F}_A(T) \rangle &\simeq \mathcal{U}^{ijpq} \left\{ k_A^4 e^{-\Gamma_A T} \nabla_{\mathbf{R}} \left[ \text{Re } G_{ij}^{(0)}(\mathbf{R}, \omega_A) \text{Re } G_{pq}^{(0)}(\mathbf{R}, \omega_A) - \text{Im } G_{ij}^{(0)}(\mathbf{R}, \omega_A) \text{Im } G_{pq}^{(0)}(\mathbf{R}, \omega_A) \right] \right. \\ &+ \frac{k_B^4 \Omega^2 e^{-(\Gamma_A + \Gamma_B)T/2}}{2(\Delta_{AB}^2 - \Omega^2)} \nabla_{\mathbf{R}} \left[ \left( \text{Re } G_{ij}^{(0)}(\mathbf{R}, \omega_B) \text{Re } G_{pq}^{(0)}(\mathbf{R}, \omega_B) - \text{Im } G_{ij}^{(0)}(\mathbf{R}, \omega_B) \text{Im } G_{pq}^{(0)}(\mathbf{R}, \omega_B) \right) \right. \\ &\times (\cos(\Delta_{AB}T) + \cos[\Delta_{AB}(T - \pi/\Omega)]) - 2 \text{Re } G_{ij}^{(0)}(\mathbf{R}, \omega_B) \text{Im } G_{pq}^{(0)}(\mathbf{R}, \omega_B) \\ &\left. \left. \times (\sin(\Delta_{AB}T) + \sin[\Delta_{AB}(T - \pi/\Omega)]) \right] \right\} \end{aligned}$$

$$\begin{aligned} \langle \mathbf{F}_B(T) \rangle &\simeq -\mathcal{U}^{ijpq} \left\{ k_A^4 e^{-\Gamma_A T} \nabla_{\mathbf{R}} \left[ \text{Re } G_{ij}^{(0)}(\mathbf{R}, \omega_A) \text{Re } G_{pq}^{(0)}(\mathbf{R}, \omega_A) + \text{Im } G_{ij}^{(0)}(\mathbf{R}, \omega_A) \text{Im } G_{pq}^{(0)}(\mathbf{R}, \omega_A) \right] \right. \\ &+ \frac{k_A^2 k_B^2 \Omega^2 e^{-(\Gamma_A + \Gamma_B)T/2}}{2(\Delta_{AB}^2 - \Omega^2)} \nabla_{\mathbf{R}} \left[ \left( \text{Re } G_{ij}^{(0)}(\mathbf{R}, \omega_A) \text{Re } G_{pq}^{(0)}(\mathbf{R}, \omega_B) + \text{Im } G_{ij}^{(0)}(\mathbf{R}, \omega_A) \text{Im } G_{pq}^{(0)}(\mathbf{R}, \omega_B) \right) \right. \\ &\times (\cos(\Delta_{AB}T) + \cos[\Delta_{AB}(T - \pi/\Omega)]) - \left( \text{Re } G_{ij}^{(0)}(\mathbf{R}, \omega_A) \text{Im } G_{pq}^{(0)}(\mathbf{R}, \omega_B) \right. \\ &\left. \left. - \text{Im } G_{ij}^{(0)}(\mathbf{R}, \omega_A) \text{Re } G_{pq}^{(0)}(\mathbf{R}, \omega_B) \right) (\sin(\Delta_{AB}T) + \sin[\Delta_{AB}(T - \pi/\Omega)]) \right] \right\} \end{aligned}$$

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# i- Computation of the dynamical vdW forces



$$\langle \mathbf{F}_A(T) \rangle \simeq \mathcal{U}^{ijpq} k_A^4 e^{-\Gamma_A T} \nabla_{\mathbf{R}} \left[ \text{Re } G_{ij}^{(0)}(\mathbf{R}, \omega_A) \text{Re } G_{pq}^{(0)}(\mathbf{R}, \omega_A) \ominus \text{Im } G_{ij}^{(0)}(\mathbf{R}, \omega_A) \text{Im } G_{pq}^{(0)}(\mathbf{R}, \omega_A) \right]$$

$$\langle \mathbf{F}_B(T) \rangle \simeq -\mathcal{U}^{ijpq} k_A^4 e^{-\Gamma_A T} \nabla_{\mathbf{R}} \left[ \text{Re } G_{ij}^{(0)}(\mathbf{R}, \omega_A) \text{Re } G_{pq}^{(0)}(\mathbf{R}, \omega_A) \oplus \text{Im } G_{ij}^{(0)}(\mathbf{R}, \omega_A) \text{Im } G_{pq}^{(0)}(\mathbf{R}, \omega_A) \right]$$

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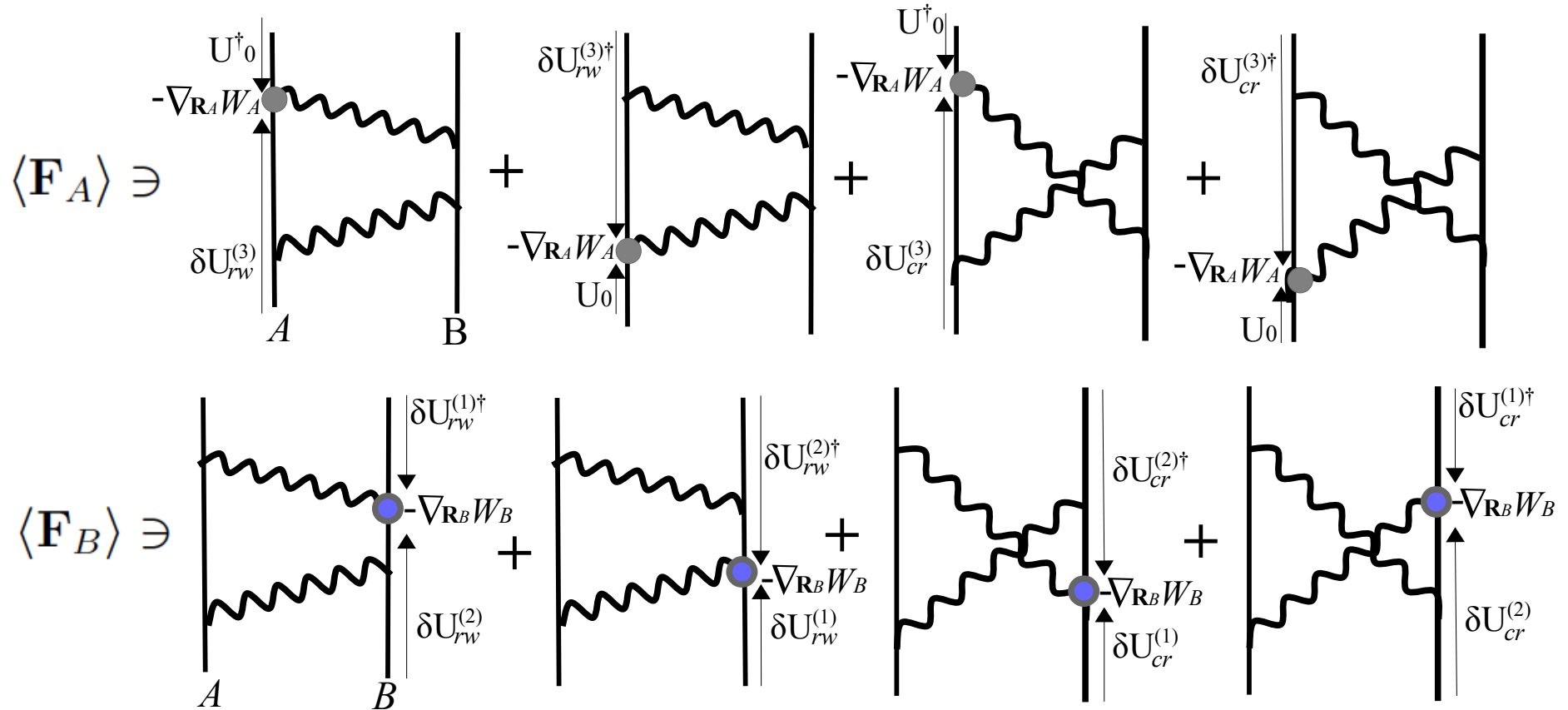
# i- Computation of the dynamical vdW forces

adiabatic excitation,

$$\Omega/\Delta_{AB} \rightarrow 0$$

Rotating wave

Counter-rotating



$$\langle \mathbf{F}_A + \mathbf{F}_B \rangle = \frac{2\omega_B \Delta_{AB}}{\omega_A^2 - \omega_B^2} \mathcal{U}^{ijpq} k_A^4 e^{-\Gamma_A T} \nabla_{\mathbf{R}} [\text{Im} G_{ij}^{(0)}(\mathbf{R}, \omega_A) \text{Im} G_{pq}^{(0)}(\mathbf{R}, \omega_A)]$$

**=> Violation of action-reaction:**

- Q1. Why the rhs is non-zero?  
Q2. Where momentum goes?

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## ii- Violation of action-reaction. Q2. Where does the momentum go?

A2. .... to the EM vacuum (total momentum conservation)

Time-variation of vacuum momentum = - Total force on the system

$$\mathbf{P} = \sum_{\mathbf{k}, \epsilon} \hbar \mathbf{k} (a_{\mathbf{k}, \epsilon}^\dagger a_{\mathbf{k}, \epsilon} + 1/2) \quad \langle \dot{\mathbf{P}} \rangle = -\langle \mathbf{F}_A + \mathbf{F}_B \rangle \quad \langle \dot{\mathbf{P}} \rangle = \langle \dot{\mathbf{P}}_{rw}^{1\gamma} \rangle + \langle \dot{\mathbf{P}}_{cr}^{1\gamma} \rangle + \langle \dot{\mathbf{P}}_{cr}^{2\gamma} \rangle$$

## ii- Violation of action-reaction. Q2. Where does momentum go?

A2. .... to the EM vacuum (total momentum conservation)

Time-variation of vacuum momentum = - Total force on the system

$$\bullet \mathbf{P} = \sum_{\mathbf{k}, \epsilon} \hbar \mathbf{k} (a_{\mathbf{k}, \epsilon}^\dagger a_{\mathbf{k}, \epsilon} + 1/2) \quad \langle \dot{\mathbf{P}} \rangle = -\langle \mathbf{F}_A + \mathbf{F}_B \rangle \quad \langle \dot{\mathbf{P}} \rangle = \langle \dot{\mathbf{P}}_{rw}^{1\gamma} \rangle + \langle \dot{\mathbf{P}}_{cr}^{1\gamma} \rangle + \langle \dot{\mathbf{P}}_{cr}^{2\gamma} \rangle$$

$$\left[ \frac{d}{dT} \left\langle \dot{\mathbf{P}} \right\rangle_{rw}^{1\gamma} \right] = \sum_{\mathbf{k}} \partial_T \langle \Psi_0 | \delta U_{rw}^{\dagger(1)}(T, \mathbf{k}) \hbar \mathbf{k} \delta U_{rw}^{(3)}(T, \mathbf{k}) | \Psi_0 \rangle + c.c. = \sum_{\mathbf{k}} \left[ \langle \Psi_0 | U_0^\dagger(T) \nabla_{\mathbf{R}_A} W_{A, \mathbf{k}} \delta U_{rw}^{(3)}(T, \mathbf{k}) | \Psi_0 \rangle + \langle \Psi_0 | \delta U_{rw}^{\dagger(1)}(T, \mathbf{k}) \nabla_{\mathbf{R}_B} W_{B, \mathbf{k}} \delta U_{rw}^{(2)}(T) | \Psi_0 \rangle \right] + c.c.$$

## ii- Violation of action-reaction. Q2. Where does momentum go?

A2. .... to the EM vacuum (total momentum conservation)

Time-variation of vacuum momentum = - Total force on the system

$$\bullet \mathbf{P} = \sum_{\mathbf{k}, \epsilon} \hbar \mathbf{k} (a_{\mathbf{k}, \epsilon}^\dagger a_{\mathbf{k}, \epsilon} + 1/2) \quad \langle \dot{\mathbf{P}} \rangle = -\langle \mathbf{F}_A + \mathbf{F}_B \rangle \quad \langle \dot{\mathbf{P}} \rangle = \langle \dot{\mathbf{P}}_{rw}^{1\gamma} \rangle + \langle \dot{\mathbf{P}}_{cr}^{1\gamma} \rangle + \langle \dot{\mathbf{P}}_{cr}^{2\gamma} \rangle$$

*rw*

$$d/dT \left\langle \left\langle \Psi_0 \left| \delta U_{rw}^{(1)\dagger} \right. \right. \right\rangle = \left\langle \left. \left. \delta U_{rw}^{(3)\dagger} \right. \right. \right\rangle + \left\langle \left. \left. \delta U_{rw}^{(2)\dagger} \right. \right. \right\rangle$$

$$\langle \dot{\mathbf{P}} \rangle_{rw}^{1\gamma} = \sum_{\mathbf{k}} \partial_T \langle \Psi_0 | \delta U_{rw}^{\dagger(1)}(T, \mathbf{k}) \hbar \mathbf{k} \delta U_{rw}^{(3)}(T, \mathbf{k}) | \Psi_0 \rangle + c.c. = \sum_{\mathbf{k}} \left[ \langle \Psi_0 | U_0^\dagger(T) \nabla_{\mathbf{R}_A} W_{A, \mathbf{k}} \delta U_{rw}^{(3)}(T, \mathbf{k}) | \Psi_0 \rangle + \langle \Psi_0 | \delta U_{rw}^{\dagger(1)}(T, \mathbf{k}) \nabla_{\mathbf{R}_B} W_{B, \mathbf{k}} \delta U_{rw}^{(2)}(T) | \Psi_0 \rangle \right] + c.c.$$

*cr*

$$d/dT \left\langle \left\langle \Psi_0 \left| \delta U_{cr}^{(1)\dagger} \right. \right. \right\rangle = \left\langle \left. \left. \delta U_{cr}^{(3)\dagger} \right. \right. \right\rangle + \left\langle \left. \left. \delta U_{cr}^{(2)\dagger} \right. \right. \right\rangle$$

$$\langle \dot{\mathbf{P}} \rangle_{cr}^{1\gamma} = \sum_{\mathbf{k}} \partial_T \langle \Psi_0 | \delta U_{cr}^{\dagger(1)}(T, \mathbf{k}) \hbar \mathbf{k} \delta U_{cr}^{(3)}(T, \mathbf{k}) | \Psi_0 \rangle + c.c. = \sum_{\mathbf{k}} \langle \Psi_0 | U_0^\dagger(T) \nabla_{\mathbf{R}_A} W_{A, \mathbf{k}} \delta U_{cr}^{(3)}(T, \mathbf{k}) | \Psi_0 \rangle + c.c.$$

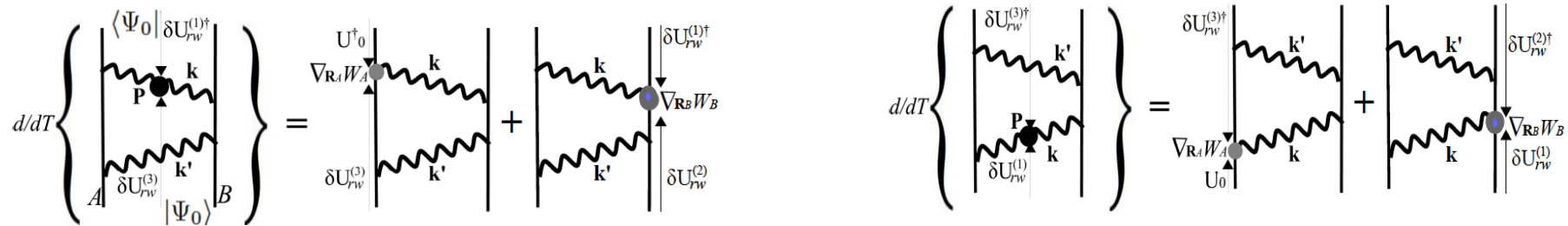
$$d/dT \left\langle \left\langle \Psi_0 \left| \delta U_{cr}^{(2)\dagger} \right. \right. \right\rangle = \left\langle \left. \left. \delta U_{cr}^{(2)\dagger} \right. \right. \right\rangle + \left\langle \left. \left. \delta U_{cr}^{(1)\dagger} \right. \right. \right\rangle$$

$$\langle \dot{\mathbf{P}} \rangle_{cr}^{2\gamma} = \sum_{\mathbf{k}, \mathbf{k}'} \partial_T \langle \Psi_0 | \delta U_{cr}^{\dagger(2)}(T, \mathbf{k}', \mathbf{k}) \hbar (\mathbf{k} + \mathbf{k}') \delta U_{cr}^{(2)}(T, \mathbf{k}, \mathbf{k}') | \Psi_0 \rangle = \sum_{\mathbf{k}, \mathbf{k}'} \langle \Psi_0 | \delta U_{cr}^{\dagger(2)}(T, \mathbf{k}', \mathbf{k}) \nabla_{\mathbf{R}_B} W_{B, \mathbf{k}} \delta U_{cr}^{(1)}(T, \mathbf{k}) | \Psi_0 \rangle + c.c.$$

## ii- Violation of action-reaction. Q1. Why does vacuum momentum not vanish?

A1. ....Optical theorem on 'will-be' emitted photons (conservation of total emission probability)

$$\bullet \mathbf{P} = \sum_{\mathbf{k}, \epsilon} \hbar \mathbf{k} (a_{\mathbf{k}, \epsilon}^\dagger a_{\mathbf{k}, \epsilon} + 1/2) \quad \langle \dot{\mathbf{P}} \rangle = -\langle \mathbf{F}_A + \mathbf{F}_B \rangle \quad \langle \dot{\mathbf{P}} \rangle = \langle \dot{\mathbf{P}}_{rw}^{1\gamma} \rangle + \langle \dot{\mathbf{P}}_{cr}^{1\gamma} \rangle + \langle \dot{\mathbf{P}}_{cr}^{2\gamma} \rangle$$



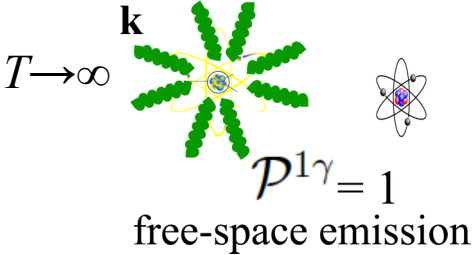
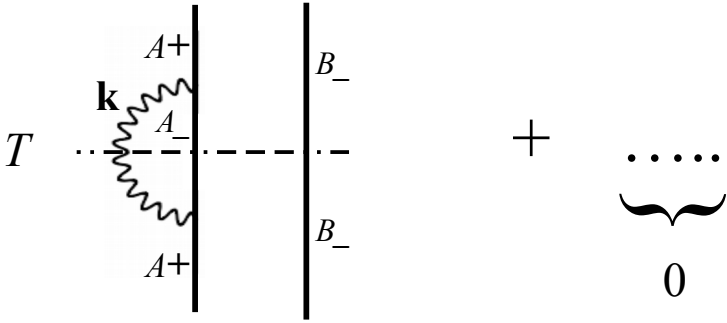
$$\langle \dot{\mathbf{P}} \rangle_{rw}^{1\gamma} = \sum_{\mathbf{k}} \hbar \mathbf{k} \partial_T \langle \Psi_0 | \delta U_{rw}^{\dagger(1)}(T, \mathbf{k}) \delta U_{rw}^{(3)}(T, \mathbf{k}) | \Psi_0 \rangle + \text{c.c.} = \sum_{\mathbf{k}} \hbar \mathbf{k} \dot{\mathcal{P}}_{rw}^{1\gamma}(\mathbf{k})$$

$$\neq 0 \Rightarrow \dot{\mathcal{P}}^{1\gamma}(\mathbf{k}) \text{ non-invariant under } \mathbf{k} \rightarrow -\mathbf{k}$$

# ii- Violation of action-reaction. Q1. Why does vacuum momentum not vanish?

One-photon emission probability  $\mathcal{P}_{rw}^{1\gamma} = 1$

P.R. Berman,  
Phys. Rev. A **76**, 043816 (2007)

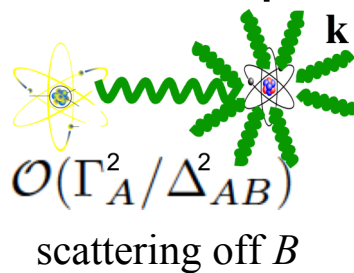
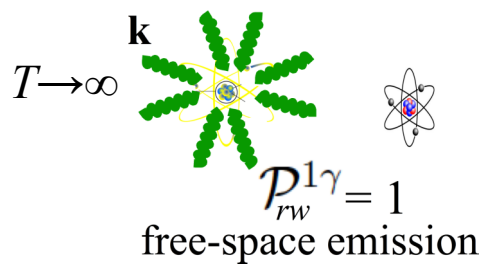
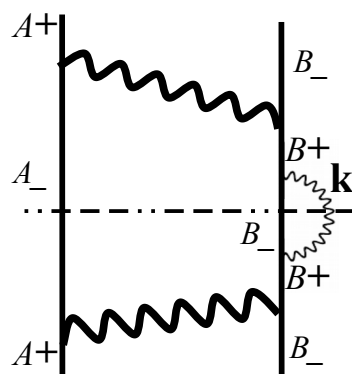
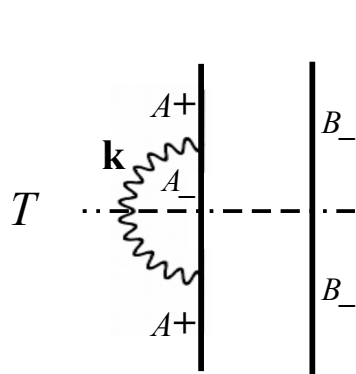




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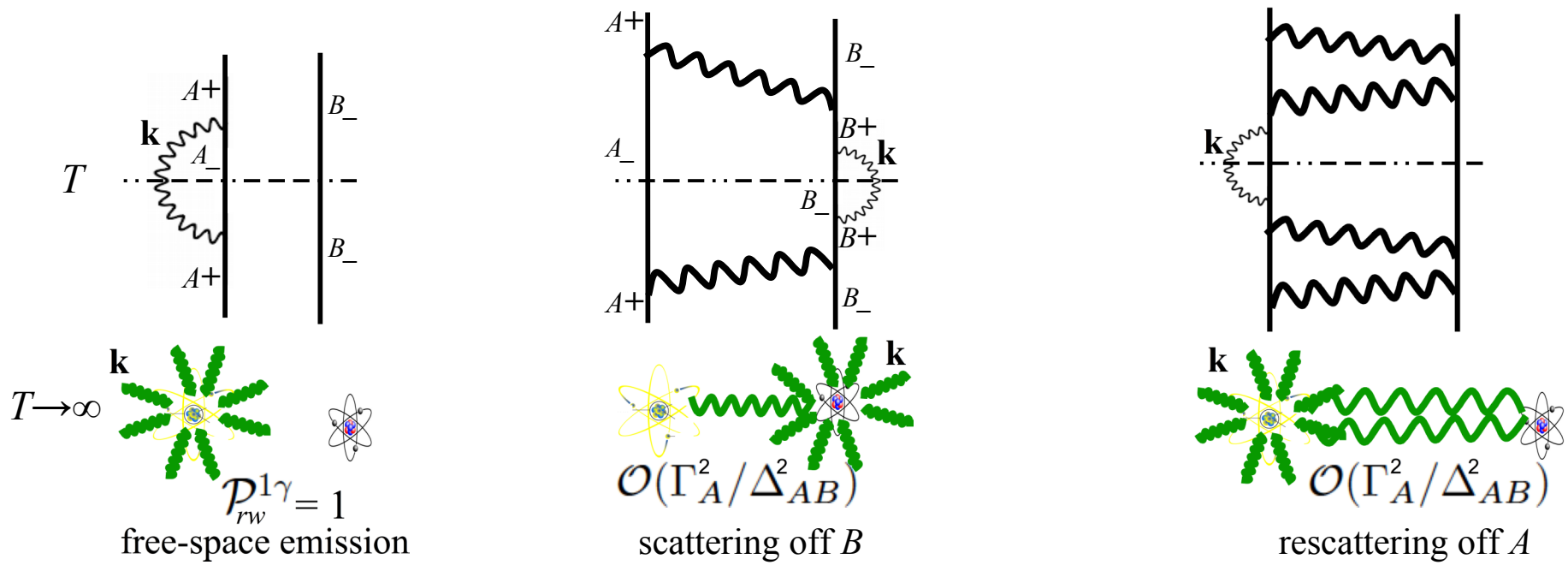
P.R. Berman,  
Phys. Rev. A **76**, 043816 (2007)



# ii- Violation of action-reaction. Q1. Why does vacuum momentum not vanish?

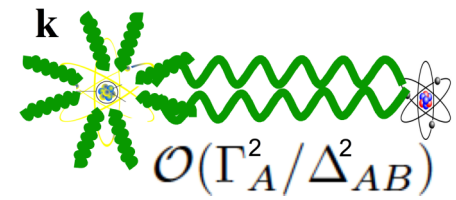
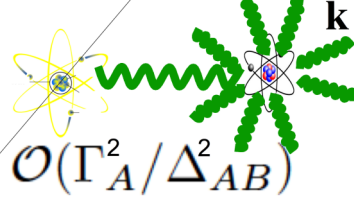
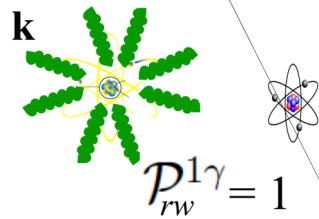
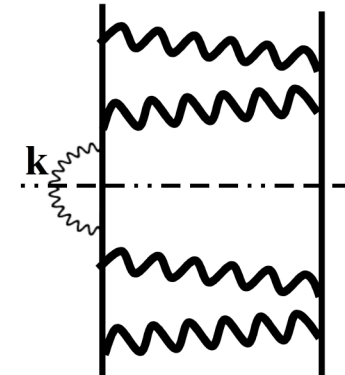
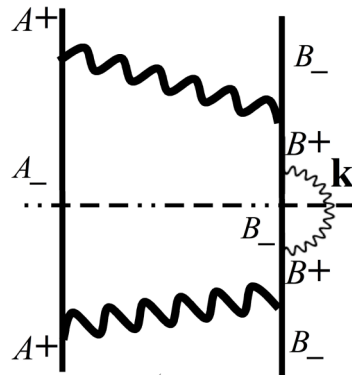
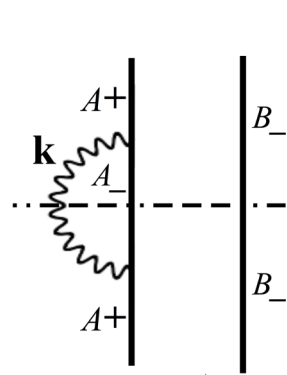
One-photon emission probability  $\mathcal{P}_{rw}^{1\gamma} = 1$

P.R. Berman,  
Phys. Rev. A **76**, 043816 (2007)

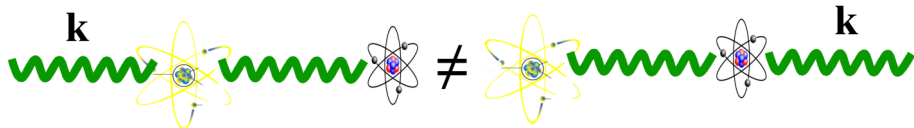
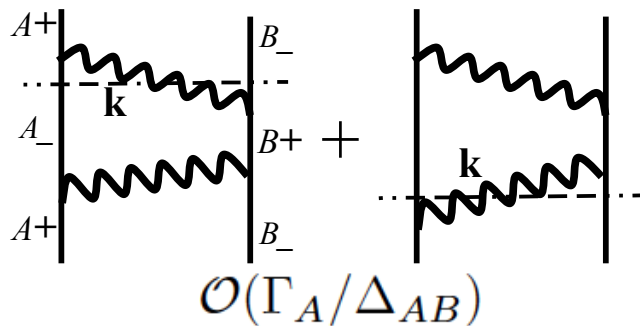


## ii- Violation of action-reaction. Q1. Why does vacuum momentum not vanish?

One-photon emission probability  $\mathcal{P}_{rw}^{1\gamma}$



Interference terms



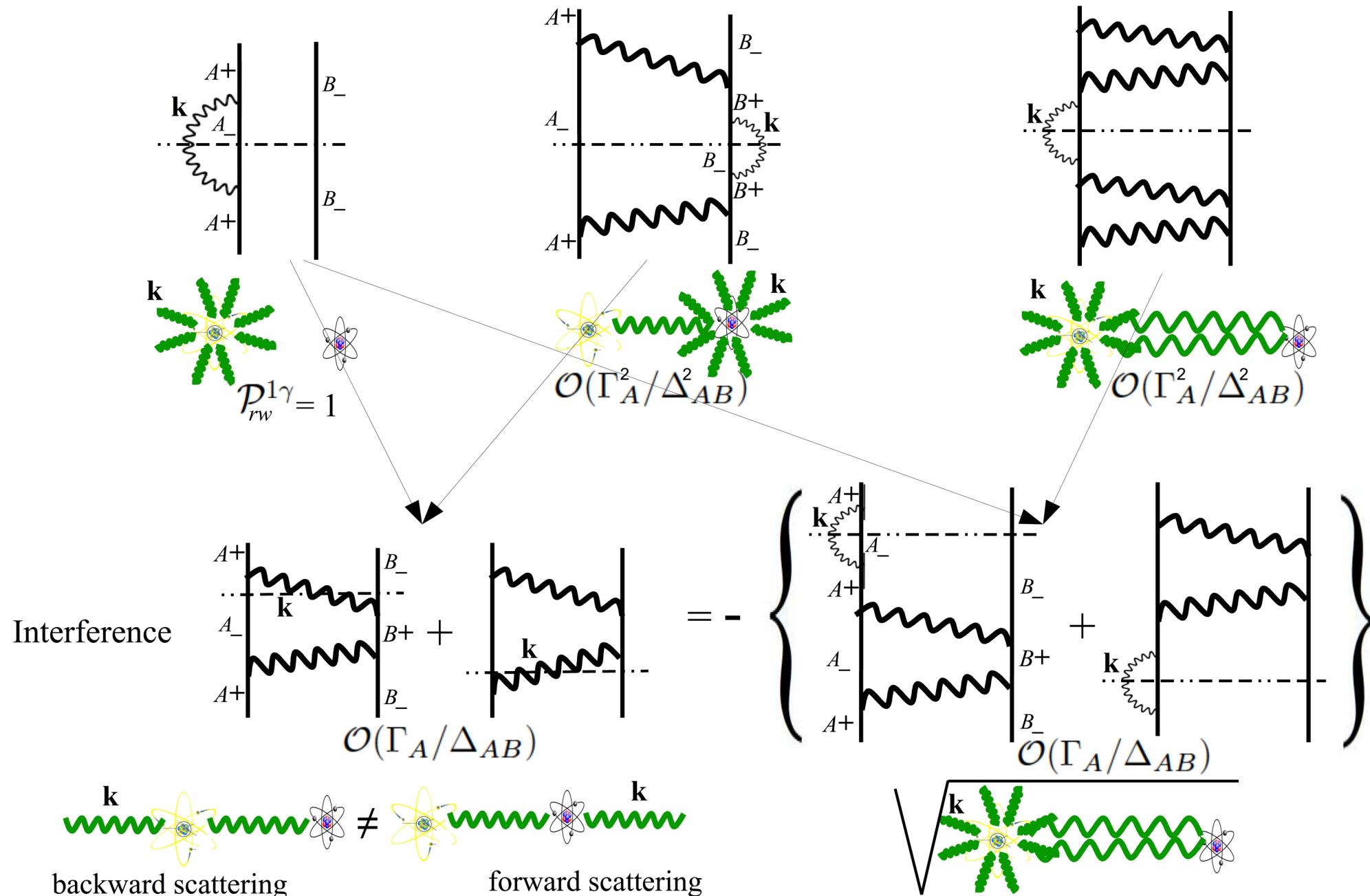
backward scattering

forward scattering

Asymmetric emission

# ii- Violation of action-reaction. Q1. Why does vacuum momentum not vanish?

One-photon emission probability  $\mathcal{P}_{rw}^{1\gamma}$  Optical theorem on 'will-be' emitted photons

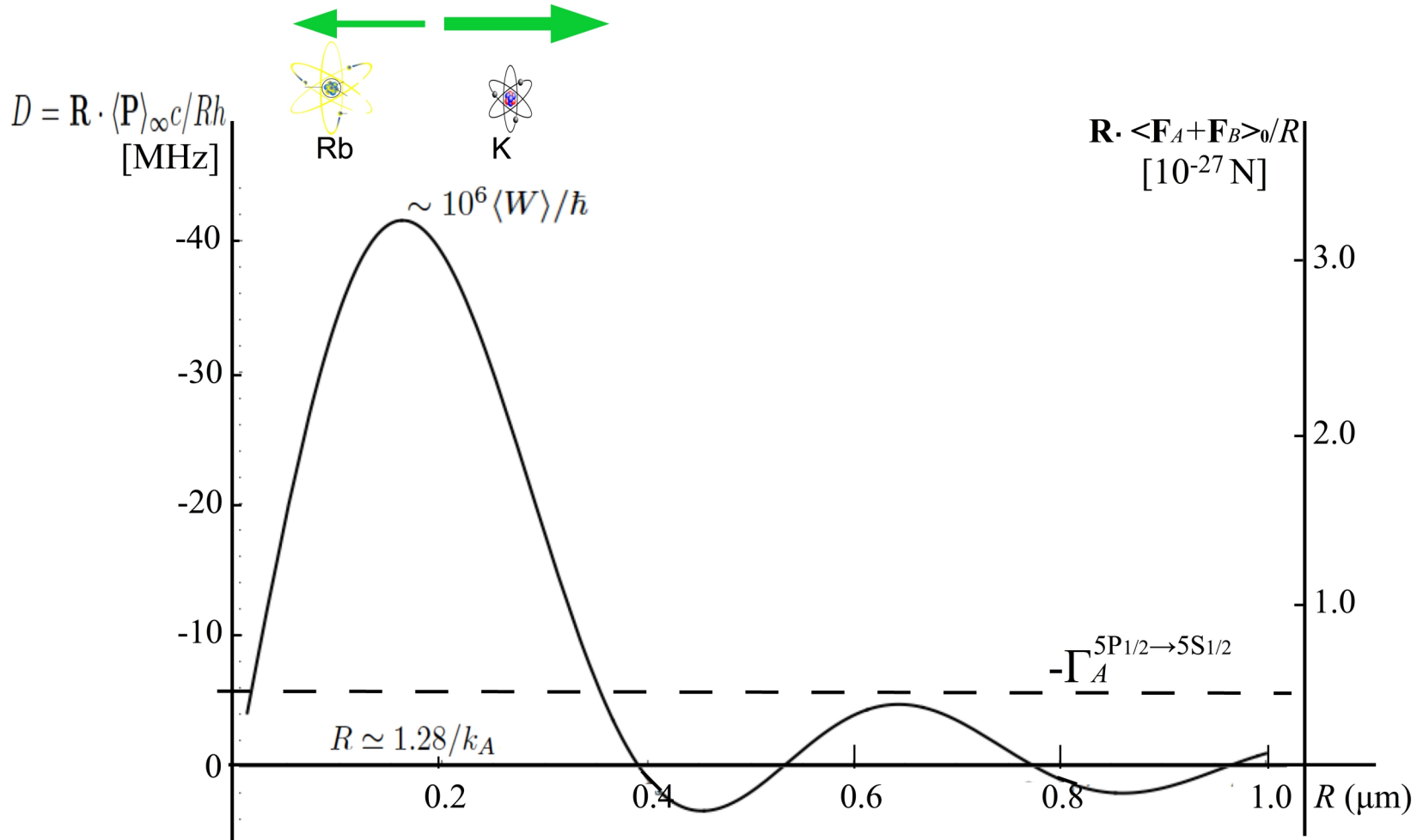


Asymmetry => **directional emission**

## ii- Violation of action-reaction. Directional spontaneous emission

Experimental detection

Action-reaction violation  $\Leftrightarrow$  Directional spontaneous emission



Rubidium atom excited at  $5P_{1/2}$  at a distance  $R$  from a potassium atom in its ground state

M. Donaire,  
arXiv:1604.07071 (2016)

M. Donaire **PSAS2016**

# Conclusions

- Violation of the action-reaction principles occurs in a two-atom system, with one atom excited.

$$(F_A + F_B)_{max} \Rightarrow D_{max} \sim \frac{d_B^2 \omega_B \omega_A^4}{(\omega_A^2 - \omega_B^2) c^3 \epsilon_0 \hbar}$$

M. Donaire,  
arXiv:1604.07071 (2016)

M. Donaire,  
Phys. Rev. A **93**, 052706 (2016)

M. Donaire, R. Guerout, A. Lambrecht,  
Phys.Rev.Lett. **115**, 033201 (2015)

- It implies **i)** the transfer of momentum to the EM vacuum and **ii)** directional spontaneous emission, which are the result of, respectively,

1) Total momentum conservation

2) The optical theorem in asymmetrically excited system  $\Leftarrow$  unitarity + locality

} common to any QFT  $\Rightarrow$

Conjecture: action-reaction violation holds in any asymmetrically excited system